

1 **Impact of microphysics on tropical precipitation**
2 **extremes in a global storm-resolving model**

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5 **Key Points:**

- 6 • Changes in microphysics influence tropical precipitation extremes in a global storm-
7 resolving model.
8 • Hourly precipitation extremes are influenced dynamically through convective up-
9 draft speed, which depends on the raindrop terminal velocity.
10 • Daily precipitation extremes are more sensitive to the microphysical modulation
11 on convective organization.

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Abstract

The impact of microphysics on tropical precipitation extremes is explored with a global storm-resolving model by modifying the terminal velocity of raindrops. Depending on the time scales, precipitation extremes respond differently. Hourly extremes are influenced dynamically through convective updraft speed, as a faster terminal velocity of raindrops increases the updraft speed by reducing the total rain in the atmosphere which increases the updraft buoyancy. However, the response of daily precipitation extremes is more sensitive to the microphysical modulation on convective organization. By being more organized with decreasing terminal velocity, daily precipitation extremes are enhanced due to increased precipitation efficiency and intensified updrafts. Thus, the results suggest that microphysics, despite often occurring at small scales, can influence the circulation at larger scales, and the microphysical imprint across different scales plays an important role in regulating tropical precipitation extremes.

Plain Language Summary

We use a global high-resolution climate model to explore the response of tropical extreme precipitation to processes governing the rain formation. We artificially alter the fall speed of raindrops to investigate its impact. The results show that it not only affects local short-duration precipitation extremes by changing the updraft speed, but also has the ability to modulate the spatial distributions of precipitation, which in the end influences precipitation extremes accumulated over longer time scales.

1 Introduction

Precipitation extremes have long been posing tremendous threats to our society, and global warming adds extra uncertainties and likely exacerbates the situation. The uncertainties come partially from the variable nature of the extremes. Unlike mean precipitation which is constrained energetically (Allen & Ingram, 2002; Held & Soden, 2006), extreme precipitation, occurring at small scales, can be sensitive to many local influences (O’Gorman, 2015) and, therefore, is much less known.

The conventional climate models generally struggle to simulate extremes, and cannot reach a consensus in terms of how the tropical daily precipitation extremes respond to climate change (O’Gorman & Schneider, 2009). The disagreement comes primarily from the inconsistencies in convective updraft speed, which is attributed to the model deficiency to represent convective processes due to the coarse model resolution (typically at 100 km) and the use of convective parameterization. Because of the model deficiencies in parameterized convection, the impact from the microphysics is often obscured.

In recent decades, with computational and technological advances, it has become feasible to run simulations at a resolution (≤ 5 km) at which convective parameterization can be switched off. Such high-resolution simulations are, however, often restricted to be over relatively small domains instead of the entire globe. Nevertheless, progress has been made in understanding tropical precipitation extremes. One important finding is the recognition of the microphysical changes in modulating precipitation extremes. Using idealized simulations configured in radiative-convective equilibrium (RCE) without rotation, Parodi and Emanuel (2009) found that the terminal velocity of raindrops determines the convective updraft speed through the condensate loading effect, which acts, dynamically, to alter precipitation extremes. With a similar setup, Singh and O’Gorman (2014) found that the response of extreme precipitation to warming depends on the choice of microphysics scheme, and that this dependency mainly comes from the effective hydrometeor fall speed simulated by different schemes which affects precipitation efficiency. Another important finding is that the response of precipitation extremes to warming may be related to changes in convective organization. Again in idealized RCE simulations with

61 homogeneous boundary conditions and no rotation, tropical convection has been shown
 62 to be able to spontaneously organize into a large convective cluster, which is referred to
 63 as convective self-aggregation (Held et al., 1993; Bretherton et al., 2005; Wing et al., 2018).
 64 The degree of organization in these studies can vary with changing setups and several
 65 studies found notable increases in precipitation extremes when convection becomes more
 66 aggregated (Bao et al., 2017; Pendergrass et al., 2016; Fildier et al., 2020). Bao and Sher-
 67 wood (2019) showed that daily precipitation extremes increase in a more organized state
 68 because organization increases the precipitation duration while instantaneous precipi-
 69 tation extremes are almost not affected. Fildier et al. (2020) found that organization in-
 70 tensifies hourly precipitation extremes by increasing precipitation efficiency. One ma-
 71 jor concern of these studies is the small domain and the simple idealization adopted in
 72 RCE. As a result, processes occurring at scales that are beyond the limit of the domain
 73 size are missing.

74 In this study, we use a realistically configured global storm-resolving model to in-
 75 vestigate the role of microphysics in tropical precipitation extremes. Similar to Parodi
 76 and Emanuel (2009), the microphysical element we focus on is the terminal velocity of
 77 raindrops. As the simulation covers a global domain, large-scale circulation and its im-
 78 pact on tropical convection and precipitation extremes are included. Additionally, we
 79 expect that the microphysical processes, despite happening at convective scales, may feed
 80 back to larger scales. The rest of the article is organized as follows: section 2 describes
 81 model details and experiments, section 3 shows results, and our discussion and conclu-
 82 sions are given in section 4.

83 2 Model and experiments

84 The simulations are conducted with ICON (Icosahedral Nonhydrostatic Weather
 85 and Climate Model; Zängl et al., 2015) at a quasi-uniform horizontal mesh of 5 km. The
 86 experiments are configured following the experimental protocol for DYAMOND (The DY-
 87 namics of the Atmospheric general circulation Modeled On Non-hydrostatic Domains;
 88 Stevens et al., 2019), in which the global meteorological analysis from the European Cen-
 89 ter for Medium Range Weather Forecasts (ECMWF) are used to initialize the model and
 90 daily observed sea surface temperatures are forced as boundary conditions. All the sim-
 91 ulations are run for 20 days from August 1st in 2016, and the hourly output for the last
 92 5 days over the tropical ocean grids (30°N-10°S) are used in the analysis.

93 The microphysics scheme (Baldauf et al., 2011) used has five hydrometeor species
 94 (rain, snow, graupel, cloud ice and cloud water). It is a single-moment scheme in which
 95 the precipitation particles are assumed to be exponentially distributed in size with re-
 96 spect to particle diameter. The terminal velocity of individual raindrops in this scheme
 97 is assumed to be only related to drop size. We change the terminal velocity of raindrops
 98 (V_{rain}) by rescaling the original formula with a fixed coefficient (Table 1), including a
 99 control simulation with the default V_{rain} (Ct), an increased velocity simulation with dou-
 100 bled V_{rain} (Db), and two decreased velocity simulations with quartered V_{rain} (Qt) and
 101 halved V_{rain} (Hf). The microphysical process is perturbed in an extreme way to inves-
 102 tigate the impact. Unlike Parodi and Emanuel (2009) who adopted a fixed velocity for
 103 all raindrops, our method perturbs the relative magnitude of the fall speed and, thus,
 104 should not substantially alter the particle interactions. The other physical parameter-
 105 izations include a radiation scheme (Rapid Radiative Transfer Model; Mlawer et al., 1997),
 106 a turbulent mixing scheme (Raschendorfer, n.d.) based on a prognostic equation for tur-
 107 bulent kinetic energy (TKE) and an interactive surface flux scheme and soil model (Schrodin
 108 & Heise, 2002). Further details about ICON and the DYAMOND configuration are given
 109 by Hohenegger et al. (2020).

Table 1. Acronyms of the experiments, the corresponding microphysical modifications to the terminal velocity of raindrops (V_{rain} * rescaling coefficient), statistics of the convective organization metric I_{org} diagnosed from precipitable water (PW) and precipitation (PR), and the mean net atmospheric radiation for all-sky ($R_{a,total}$: W m^{-2}) and clear-sky ($R_{a,cs}$: W m^{-2}) conditions.

Name	Rescaling coefficient	$I_{org}(\text{PW})$	$I_{org}(\text{PR})$	$R_{a,total}$	$R_{a,cs}$
Qt	0.25	0.939	0.874	-98.3	-118.9
Hf	0.5	0.861	0.841	-105.2	-119.5
Ct	1.0	0.834	0.803	-109.0	-119.8
Db	2.0	0.829	0.806	-111.5	-120.7

3 Results

3.1 Thermodynamic characteristics of the tropical mean state

We first focus on some of the thermodynamic characteristics of the tropical mean state. Figure 1a and 1b show the differences in the virtual temperature (T_v) and relative humidity (RH) of the runs with modified terminal velocity relative to the control run (Ct). With a slower terminal velocity of raindrops, the troposphere becomes more stable as the free troposphere is substantially warmer, whereas the boundary layer is colder. Meanwhile, the entire troposphere becomes more humid especially between 600hPa and 800hPa. A cooler and moister boundary layer can be attributed to the slower raindrop velocity which increases the residence time of the raindrops and enhances evaporation (Fig. 1c). A warmer free troposphere can be explained from the changes in humidity, as the troposphere gets moister, it better protects rising convective parcels from the impact of entrainment, which, as a result, ensures a more precise moist-adiabatic ascent (Singh & O’Gorman, 2013; Seeley & Romps, 2015). Gravity waves then act to quickly adjust the temperature in the non-convective regions and homogenize the temperature horizontally in the free troposphere (Bretherton & Smolarkiewicz, 1989; Sobel & Bretherton, 2000). Thus, the tropical troposphere becomes warmer and more stable with a decreased raindrop velocity. On the other hand, an increased raindrop terminal velocity leads to the opposite response by reducing the residence time of the raindrops and suppressing evaporation, promoting a colder and less stable troposphere. Thus, the change in the microphysics which happens at small scales is shown to impact the tropics as a whole.

3.2 Tropical precipitation extremes

Fractional changes of precipitation as a function of precipitation percentile are shown in Figure 1d. Here we compare hourly and daily precipitation calculated from simulations with different raindrop terminal velocities relative to Ct. Extreme precipitation exhibits distinct variations at different time scales. At hourly time scales, extreme precipitation increases roughly linearly with the terminal velocity. At high precipitation percentiles ($> 99.9\text{th}$), precipitation reduces by 50% in Qt and 20% in Hf while increasing by $\sim 20\%$ in Db. On the other hand, daily precipitation extremes do not seem to vary linearly with the terminal velocity. The highest daily extremes occur in the case with the slowest terminal velocity (Qt), increasing by $\sim 40\%$ at 99.99th percentile. A dramatic change in extreme precipitation from hourly to daily time scales (especially in Qt) suggests they are potentially controlled by different mechanisms.

To better understand the mechanisms controlling extreme precipitation, we apply a scaling analysis method that separates extreme precipitation into thermodynamic, dynamic and precipitation efficiency components following O’Gorman and Schneider (2009)

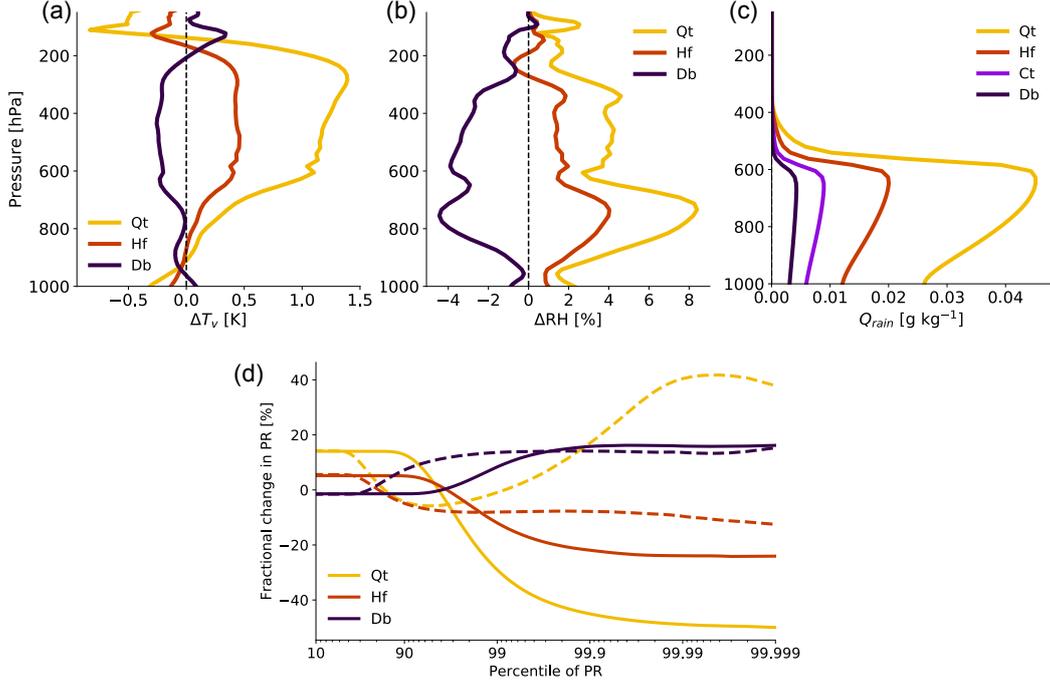


Figure 1. (a,b) Difference in tropical mean virtual temperature (T_v) and relative humidity (RH) relative to Ct. (c) Profiles of the tropical mean rain mixing ratio (Q_{rain}). (d) Fractional changes in precipitation (PR) relative to Ct as a function of PR percentile. Solid (dashed) lines represent hourly (daily) precipitation.

and Muller et al. (2011). First, a high-percentile precipitation rate (P_e) is represented by the product of net condensation rate (C) and precipitation efficiency (ϵ) conditioned on P_e :

$$P_e = \epsilon C. \quad (1)$$

The condensation rate can be approximated as:

$$C \approx \int_{p_s}^{p_t} \frac{\omega}{g} \frac{dq_s}{dp} \Big|_{\theta_e^*} dp, \quad (2)$$

where p_s and p_t are pressure at the surface and the tropopause, g is the acceleration from gravity, ω is the updraft velocity in pressure coordinates conditioned on precipitation extremes, and $\frac{dq_s}{dp} \Big|_{\theta_e^*}$ is the change in saturation specific humidity with respect to pressure at constant saturation equivalent potential temperature θ_e^* and is referred to as the moisture lapse rate. Eq. (2) assumes that condensation occurs roughly moist-adiabatically during an extreme precipitation event. The condensation here represents the net condensation (condensation minus evaporation). As a result, ϵ , differs from a conventional precipitation efficiency, is defined as the extreme precipitation rate divided by the net condensation rate. Following the steps in Fildier et al. (2020) (except using pressure coordinates here), we define a dynamical term (M) which represents the column-integrated mass flux as:

$$M = - \int_{p_s}^{p_t} \frac{\omega}{g} dp. \quad (3)$$

164 Then a thermodynamic term (Γ_q) dominated by the moisture lapse rate can be obtained
 165 by dividing the condensation rate by the dynamical term using Eq. (2),(3):

$$166 \quad \Gamma_q \approx -\frac{1}{M} \int_{p_s}^{p_t} \frac{\omega}{g} \frac{dq_s}{dp} \Big|_{\theta_s^*} dp = -\int_{p_s}^{p_t} \frac{1}{g} \frac{\omega}{M} \frac{dq_s}{dp} \Big|_{\theta_s^*} dp. \quad (4)$$

167 Note that M is a single value and can be put inside the integral to separate the dynamical
 168 component from ω . By combining Eq. (1),(2),(3),(4), we have:

$$169 \quad P_e \approx \epsilon M \Gamma_q. \quad (5)$$

170 Thus, changes in extreme precipitation can be decomposed into a dynamical component
 171 ($\frac{\delta M}{M}$) through changes in updraft velocity, a thermodynamic component ($\frac{\delta \Gamma_q}{\Gamma_q}$) through
 172 changes in the moisture lapse rate, and a precipitation efficiency component ($\frac{\delta \epsilon}{\epsilon}$):

$$173 \quad \frac{\delta P_e}{P_e} \approx \frac{\delta \epsilon}{\epsilon} + \frac{\delta M}{M} + \frac{\delta \Gamma_q}{\Gamma_q} \quad (6)$$

174 **3.2.1 Hourly precipitation extremes**

175 At hourly time scales, changes in precipitation extremes are almost entirely due
 176 to the dynamical component while the thermodynamic and efficiency components play
 177 little role (Fig. 2). As the dynamical component is controlled by updraft velocity, it indicates
 178 that the convective updraft is stronger (weaker) when the terminal velocity is faster
 179 (slower). This is confirmed in Fig. 2g showing that the updraft velocity increases following
 180 the raindrop terminal velocity throughout the entire troposphere. One potential explanation
 181 for the increase in the updraft velocity is the reduced tropospheric stability,
 182 as section 3.1 shows that the troposphere becomes increasingly unstable with a faster
 183 raindrop terminal velocity. Another possibility is the weakened condensate loading effect
 184 resulting from more rapid removal of condensates from the atmosphere when raindrops
 185 are allowed to fall faster (Parodi & Emanuel, 2009). This tends to moderate the
 186 condensate loading effect in reducing the updraft buoyancy.

187 To separate these two effects and thus to understand what controls the updraft speed
 188 when hourly precipitation extremes occur, we investigate its relationship with buoyancy
 189 (B) as the vertically integrated buoyancy provides the kinetic energy for the convective
 190 updraft:

$$191 \quad \frac{1}{2} w_{max}^2 \sim \int_0^{z_t} B dz, \quad (7)$$

192 where w_{max} is the maximum value of vertical velocity in the vertical column and z_t is
 193 fixed at about 11 km. The buoyancy of an updraft air parcel is formulated as:

$$194 \quad B \approx g \frac{T'_v}{T_v} - gl', \quad (8)$$

195 where T'_v is the virtual temperature excess between the updraft air parcel and its environment
 196 and l' is the rain water mixing ratio in the updraft column. The updraft grid
 197 cells are identified as the grid cells in which w_{max} exceeds the 99.9th percentile value.
 198 Then for each updraft, the corresponding environment is defined as the non-cloudy grid
 199 cells (the mixing ratio of the total condensates $< 10^{-5}$ kg kg $^{-1}$) within 30 km radius of
 200 the updraft grid cell. The first term of on the right-hand side in Eq. (8) is associated with
 201 the static stability of the updraft environment, and the second term depicts the reduction
 202 of the updraft buoyancy from the condensate loading as the condensates increase
 203 the effective density of the updraft system. We can test the impact of condensate loading
 204 by comparing the relationship between the updraft kinetic energy and the vertical
 205 integration of the buoyancy with the buoyancy computed either including or omitting
 206 the loading term.

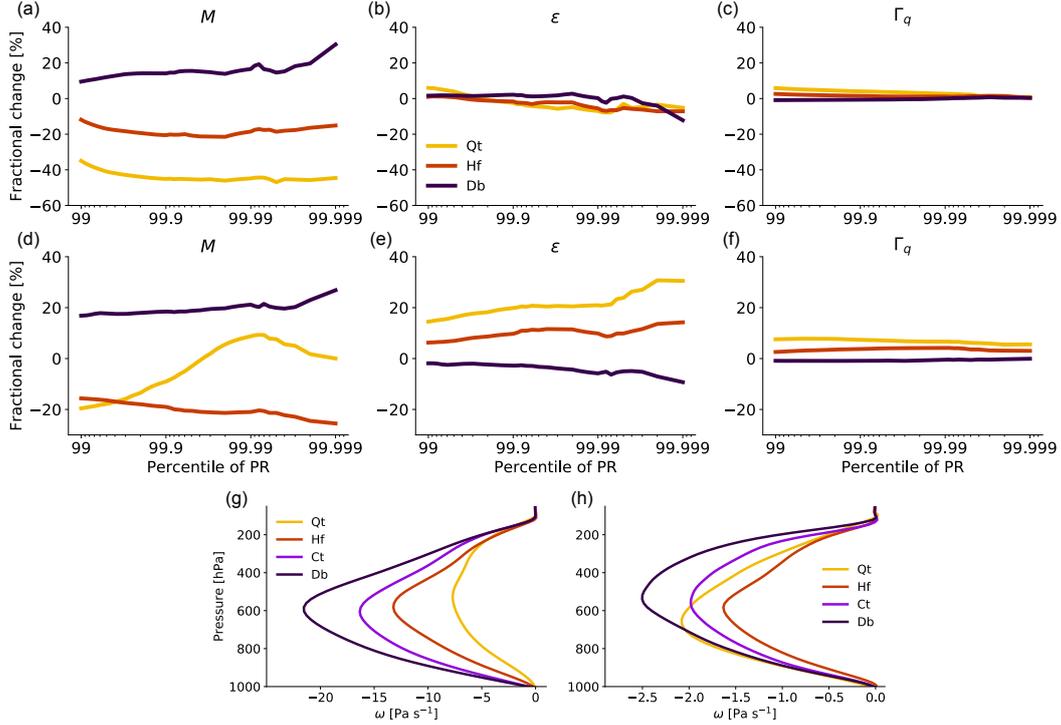


Figure 2. (Upper two rows) Fractional changes in the dynamical component (M), precipitation efficiency component (ϵ) and thermodynamic component (Γ_q) relative to Ct as a function of PR percentile. Results are shown for hourly (a-c) and daily (d-f) precipitation. (Lower row) Profiles of pressure velocity (ω) composited by 99.999th percentile of hourly (g) and daily (h) precipitation.

207 In Figure 3, we compare the updraft peak kinetic energy with the vertical integration of the buoyancy one excluding (Fig. 3a) and one including the condensate loading term (Fig. 3b). By including the condensate loading term, the updraft peak kinetic energy is much better correlated with the buoyancy integral, with the Pearson correlation coefficient (R) increasing from 0.46 to 0.73. This implies that the change in the condensate loading caused by varying raindrop terminal velocity is crucial to modulating the updraft buoyancy and thus also the updraft speed. Therefore, hourly precipitation extremes are mostly determined by the dynamical contribution from changes in the convective updraft, and the updraft speed is sensitive to microphysics through the condensate loading effect, in agreement with the results of Parodi and Emanuel (2009).

217 3.2.2 Daily precipitation extremes

218 For daily precipitation extremes, the dynamical component still plays a very im-
 219 portant role as it contributes to the highest percentiles (99.999th) roughly 20% increase
 220 in Db and 20% reduction in Hf (Fig. 2d). However, in Qt where the dramatic rise in pre-
 221 cipitation extremes occur, contributions from both the dynamics and efficiency are im-
 222 portant (Fig. 2e). While the precipitation efficiency always favors intensifying extremes
 223 in Qt, the dynamical component shifts from negative contributions at less extreme (<99.95th)
 224 percentiles to positive contributions at more extreme percentiles. Finally, the thermo-
 225 dynamic component is mainly associated with the slight temperature changes, and con-

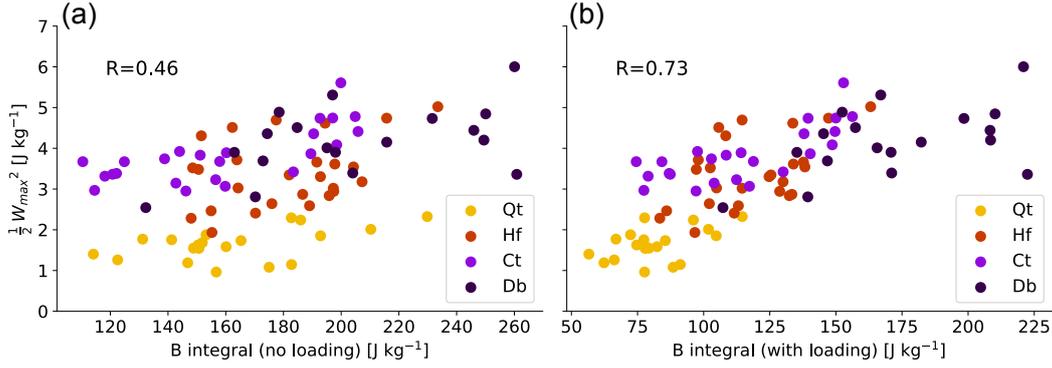


Figure 3. Peak kinetic energy ($\frac{1}{2}w_{max}^2$) vs. vertical integrated buoyancy (B) calculated without the loading effect (a) and with the loading effect (b).

226 sistent with its role at hourly time scales, its contribution to daily extremes is also less
 227 important than the other two components (Fig. 2f).

228 Figure 2e also shows that precipitation efficiency is negatively correlated with the
 229 raindrop terminal velocity. For the highest (99.999th) percentile, efficiency increases by
 230 $\sim 30\%$ in Qt and $\sim 10\%$ in Hf while reducing by $\sim 10\%$ in Db. In idealized RCE simu-
 231 lations, higher precipitation efficiency often occurs in the state of more organized con-
 232 vection (Bao & Sherwood, 2019; Fildier et al., 2020). We believe this conclusion also holds
 233 here and that the degree of convective organization plays a role in the daily precipita-
 234 tion extremes. To explore the impact of organization, first we plot the probability den-
 235 sity functions of RH . Figure 4 shows bimodal structures of the RH distributions with
 236 distinct moist and dry peaks in all cases. The bimodality is most pronounced in Qt, and
 237 becomes less pronounced with increasing raindrop velocity. As an increased moisture vari-
 238 ance is a typical feature of more organized convection (Wing et al., 2018), Fig. 4 sug-
 239 gests that convection is the most organized in Qt. To complement this qualitative inter-
 240 pretation, we quantify the degree of organization with the organization index I_{org} (Ta-
 241 ble 1), which is based on the distribution of the nearest neighbor distance between iden-
 242 tified convective clusters (Tompkins & Semie, 2017). We identify the grid cells where the
 243 precipitable water (PW) exceeds the 99.9th percentile value as convective grid cells. Two
 244 convective grid cells belong to one convective cluster if they share one boundary. The
 245 calculation of I_{org} using PW is referred to as $I_{org}(PW)$. To test the robustness of the
 246 results, we also apply the same calculation, but using daily precipitation data to iden-
 247 tify convective clusters $I_{org}(PR)$. As a higher I_{org} indicates more organized convec-
 248 tion, the result, consistent with Fig. 4, reveals that Qt, which has the highest precipitation
 249 efficiency, is the most organized, and the degree of organization decreases with increas-
 250 ing raindrop terminal velocity. At daily time scales, precipitation efficiency is mainly af-
 251 fected by the horizontal advection of water (Muller et al., 2011). In a more organized
 252 state, more moisture is transported to the column where extreme precipitation occurs.
 253 In addition to increasing precipitation efficiency, more organized convection can induce
 254 precipitation events to last longer, exerting positive dynamical contribution. This is es-
 255 pecially the case of Qt in which multiple tropical cyclones develop. Despite having the
 256 weakest updraft speed at hourly time scales, organization intensifies the mean updraft
 257 averaged over a day by being more persistent in the same locations.

258 We speculate that the varying degrees of organization here is related to the changes
 259 in the atmospheric cloud radiative effect (ACRE) and the surface enthalpy fluxes. For
 260 ACRE, as discussed in section 3.1, when the terminal velocity is greater, the troposphere
 261 becomes drier. Thus, the clouds tend to shrink and the cloud radiative effect is weak-

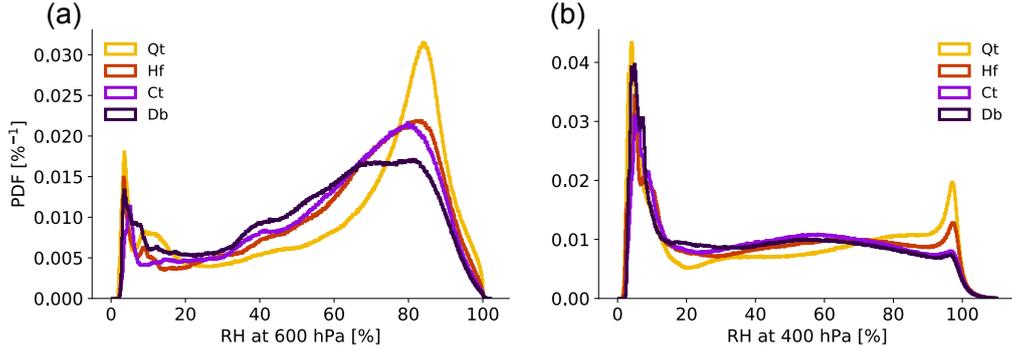


Figure 4. Probability density functions (PDF) of relative humidity (RH) distribution at 600 hPa (a) and 400 hPa (b).

262 ened. Indeed, the net atmospheric radiation under the clear-sky conditions changes little
 263 among the simulations, while it differs substantially under the all-sky conditions (Ta-
 264 ble 1), contributed mostly by the changes in the outgoing longwave radiation. As the
 265 difference in the net atmospheric radiation between the all-sky and the clear-sky con-
 266 ditions represents the cloud radiative effect, it implies that the strength of ACRE reduces
 267 with increasing terminal velocity. Such differences in ACRE, driven by the changes in
 268 the microphysics, develop very rapidly during the first few days of the simulations. They
 269 influence the net atmospheric energy uptake which, as a result, modulate the mean cir-
 270 culation and organization. This result is consistent with Popp and Bony (2019) who also
 271 found that the cloud radiative effect affects the zonal convective clustering in observa-
 272 tion. In terms of the surface fluxes, when the terminal velocity is reduced, increased evap-
 273 oration favors stronger cold pools which lead to enhanced surface fluxes. This, through
 274 a wind-induced surface heat exchange (WISHE) feedback, also contributes to a more or-
 275 ganized state. Such a feedback is especially important in Qt where multiple tropical cy-
 276 clones develop, as the feedback between surface wind and surface enthalpy flux are very
 277 fundamental for tropical cyclones (Zhang & Emanuel, 2016; Muller & Romps, 2018).

278 In contrast to hourly precipitation extremes that are mainly determined by the convective-
 279 scale dynamics, microphysics affects daily extremes by changing the behavior of organ-
 280 ization over larger scales. A more organized state primarily increases the precipitation
 281 efficiency while at the same time can lead to a more positive dynamical contribution by
 282 intensifying the updraft speed.

283 4 Discussion and conclusions

284 We use a global storm-resolving model to investigate the impact of microphysics
 285 (terminal velocity of raindrops) on tropical precipitation extremes. We find that the mi-
 286 crophysics influences hourly precipitation extremes by changing the convective updraft
 287 speed, which is fundamentally linked to the condensate loading effect dictated by the rain-
 288 drop terminal velocity in the microphysics parameterization. Contrarily, daily precipi-
 289 tation extremes are related to the microphysical influence on convective organization,
 290 as organization enhances daily precipitation extremes by higher precipitation efficiency
 291 and intensified updrafts.

292 This work highlights the importance of microphysics in tropical precipitation ex-
 293 tremes over different time scales. First it shows the dependence of convective updraft
 294 speed on raindrop terminal velocity, emphasizing the often overlooked microphysical mod-

295 ulation on convective-scale dynamics as first proposed by Parodi and Emanuel (2009).
 296 Further, it demonstrates that changes in small-scale microphysics can influence the mean
 297 climate as a whole. In particular, the microphysical imprints on convective organization
 298 can modulate precipitation extremes accumulated over long time scales, confirming the
 299 results from idealized RCE simulations that daily precipitation extremes increase when
 300 convection becomes more organized (Bao & Sherwood, 2019).

301 This work confirms, in a realistic simulation, the possibility of having varying de-
 302 grees of convective organization, as is often reported in idealized simulations of RCE. Typ-
 303 ical impacts of having convection being more organized, such as increased moisture vari-
 304 ances and enhanced precipitation efficiency, are in line with the results obtained from
 305 those RCE simulations. The different behavior of convective organization is linked to the
 306 microphysical modulation on moisture, which influences the cloud radiative effect and
 307 the surface flux feedbacks. This supports the conclusions from idealized RCE simulations
 308 that the radiative feedbacks and surface flux feedbacks are important for convective or-
 309 ganization (Wing et al., 2018).

310 The fundamental role of microphysics lies in its modification of moisture. Although
 311 the microphysical element explored in this work is the terminal velocity of raindrops, many
 312 other microphysical parameters can lead to similar changes. Thus, the results in this work
 313 should serve as an example to illustrate the non-negligible role of microphysical processes
 314 in affecting the tropical climate over a range of scales. To better simulate the tropical
 315 climate especially the precipitation extremes, an improved understanding of its interac-
 316 tion with the microphysics is hence desired.

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